LIS - CANA seminar (Marseille)

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# On the emergence of preferred structures in quantum theory

Joint work with Guilherme Franzmann & Andrea Di Biagio (soon on arXiv)

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  - Hilbert space fundamentalism (HSF) [Carroll & Singh, 2019]

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Yes, *via* the K-locality requirement

No, a structure can never be uniquely determined by H nor by  $(H, |\psi>)$ 

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Solving the tension has deep implications for the notion of emergence in QM (in physics?)

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## SUMMARY

#### Le Cotler et al.'s theorem

- 1. K-locality
- 2. K-duality
- 3. The theorem

#### II. Stoica's theorem

- 1. Core idea
- 2. K-locality is unitary-invariant
- 3. H has more symmetries than the TPS
- 4. The time-evolution symmetry

#### III. How structures acquire their identity

- 1. Which notion of uniqueness to keep?
- 2. Insights from invariant theory
- 3. Can (H,  $|\psi\rangle$ ) uniquely determine a TPS?

# L COTLER ET AL.'S THEOREM

#### I.1. K-locality

In all the talk, we fix integers n and  $(d_i)_{1 \le i \le n}$  with  $d_i \ge 2$  and a Hilbert space  $\mathcal H$  of finite dimension  $\prod_i d_i$ .

**Definition** (Tensor product structure). A TPS of  $\mathcal{H}$  is an equivalence class of isomorphisms  $T: \mathcal{H} \to \bigotimes_{i=1}^n \mathcal{H}_i$  that factorize  $\mathcal{H}$  into n factors  $\mathcal{H}_i$  of respective dimensions  $d_i$ , where two isomorphisms  $T_1$  and  $T_2$  are said to be equivalent (denoted  $T_1 \sim T_2$ ) if  $T_1T_2^{-1}$  is a product of local unitaries  $U_1 \otimes \cdots \otimes U_n$  and arbitrary permutations of the factors.

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H can be decomposed in a TPS:

Given, for each i, a choice  $(O_i^{\alpha})_{0 \leq \alpha \leq d_i^2 - 1}$  of an orthonormal basis for  $\mathcal{L}(\mathcal{H}_i)$  with  $O_i^0 = \mathbb{1}$ , any Hermitian operator  $\hat{H}$  can be uniquely decomposed as:

$$\hat{H} = a_0 \mathbb{1} + \sum_{i=1}^n \sum_{\alpha \neq 0} a_i^{\alpha} O_i^{\alpha} + \sum_{1 \leqslant i < j \leqslant n} \sum_{\alpha, \beta \neq 0} a_{ij}^{\alpha\beta} O_i^{\alpha} O_j^{\beta} + \sum_{1 \leqslant i < j < k \leqslant n} \sum_{\alpha, \beta, \gamma \neq 0} a_{ijk}^{\alpha\beta\gamma} O_i^{\alpha} O_j^{\beta} O_k^{\gamma} + \dots, \quad (1)$$

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Hence the following notion:

**Definition** (K-locality). Let  $\mathcal{T}$  be a TPS of  $\mathcal{H}$  into  $\bigotimes_{i=1}^n \mathcal{H}_i$ . For  $K \in \{1, ..., n\}$ , we say that the pair  $(\hat{H}, \mathcal{T})$  is K-local if there exists a choice of orthonormal bases  $(O_i^{\alpha})_{i,\alpha}$  for which the above decomposition (1) involves only products of at most K non-trivial operators.

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  - ▶ Ising nearest-neighbour + a term  $\sigma_z \otimes \sigma_{z+1000}$
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  - onversely, a QFT is always spacetime local, no matter K
- NB: most Hamiltonian are approx. 2-local in some TPS [Loizeau et al., 2023]

N. Loizeau, F. Morone, and D. Sels. "Unveiling order from chaos by approximate 2-localization of random matrices", Proceedings of the National Academy of Sciences (2023)

For Cotler et al., uniqueness of structures is understood **up to a global unitary** 

**Definition** (Global unitary equivalence). Two pairs  $(\hat{H}, \mathcal{T})$  and  $(\hat{H}', \mathcal{T}')$  are said equivalent if there exists a unitary  $U \in \mathcal{U}(\mathcal{H})$  such that  $(\hat{H}', \mathcal{T}') = U \cdot (\hat{H}, \mathcal{T})$ .

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**Definition** (Equivalence of TPSs). We say that  $\mathcal{T}$  and  $\mathcal{T}'$  are equivalent with respect to  $\hat{H}$  if  $T\hat{H}T^{-1}$  and  $T'\hat{H}T'^{-1}$  on  $\bigotimes_{i=1}^{n} \mathcal{H}_i$  are the same up to single-site unitaries and permutations of factors.

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• Whether H can uniquely determine a TPS can be phrased in terms of **dual TPSs** 

**Definition** (K-duality). We say that two TPSs are K-duals if they are not equivalent with respect to  $\hat{H}$  but  $\hat{H}$  is K-local in both TPSs.

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• Example of dual TPSs: Jordan-Wigner transform on the 1D Ising model

$$\mu_z^{(i)} = \prod_{j \leqslant i} \sigma_x^{(j)}$$

$$\mu_x^{(i)} = \sigma_z^{(i)} \sigma_z^{(i+1)}$$

$$\mu_x^{(n)} = \mu_z^{(n)}$$



induces

$$\hat{H} = \hat{H} = J \sum_{i} \sigma_{z}^{(i)} \sigma_{z}^{(i+1)} + h \sum_{i} \sigma_{x}^{(i)}$$

$$\downarrow$$

$$\hat{H} = J \sum_{i} \mu_{x}^{(i)} + h \sum_{i} \mu_{z}^{(i)} \mu_{z}^{(i+1)} - J \mu_{x}^{(n)} + h \mu_{z}^{(1)}$$

#### I.3. The theorem

• After a terribly intricate proof:

**Theorem** (Cotler et al.). Assume  $d_1 = \cdots = d_n \equiv d$ , and suppose the existence of a single Hamiltonian on  $\mathcal{H}$  admitting a TPS  $\mathcal{H} = \bigotimes_{i=1}^n \mathcal{H}_i$  that makes it K-local but without any K-duals. Then, if K is sufficiently small, almost all K-local Hamiltonians on  $\mathcal{H}$  do not have K-dual TPSs.

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Numerical simulations to get convinced of the validity of the assumption.

# II. STOICA'S THEOREM

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Statement.

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- Stoica understands uniqueness in a strict sense (not up to a unitary)
  - 'distinct' ≠ 'inequivalent'
- Directly follows from

**Lemma 1.** Let  $U \in \mathcal{U}(\mathcal{H})$ . The pair  $(\hat{H}, \mathcal{T})$  is K-local if and only if the pair  $U \cdot (\hat{H}, \mathcal{T}) \equiv (U\hat{H}U^{\dagger}, U \cdot \mathcal{T})$  is K-local.

**Lemma 2.** For any K-local pair  $(\hat{H}, \mathcal{T})$ , there exists at least one unitary U such that  $U\hat{H}U^{\dagger} = \hat{H}$  but  $U \cdot \mathcal{T} \neq \mathcal{T}$ .

## II.2. K-locality is unitary-invariant

Recall decomposition (1)

$$\hat{H} = a_0 \mathbb{1} + \sum_{i=1}^n \sum_{\alpha \neq 0} a_i^{\alpha} O_i^{\alpha} + \sum_{1 \leqslant i < j \leqslant n} \sum_{\alpha, \beta \neq 0} a_{ij}^{\alpha\beta} O_i^{\alpha} O_j^{\beta} + \sum_{1 \leqslant i < j < k \leqslant n} \sum_{\alpha, \beta, \gamma \neq 0} a_{ijk}^{\alpha\beta\gamma} O_i^{\alpha} O_j^{\beta} O_k^{\gamma} + \dots, \quad (1)$$

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And observe that

$$a_{ijk...}^{\alpha\beta\gamma...} = \langle \hat{H}, O_i^{\alpha} O_j^{\beta} O_k^{\gamma} \cdots \rangle_{HS}$$

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UHU† decomposed in U  $\cdot$   $\mathcal{T}$  for the bases (O<sub>i</sub> $^{\alpha}$ )<sub>i, $\alpha$ </sub> has the same decomposition as H decomposed in  $\mathcal{T}$  for the bases (U O<sub>i</sub> $^{\alpha}$ U†)<sub>i, $\alpha$ </sub>

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**Proposition.** There are symmetries of  $\hat{H}$  that are not symmetries of  $\mathcal{T}$ , that is:  $\operatorname{Stab}(\hat{H}) \not\subset \operatorname{Stab}(\mathcal{T})$ .

Short proof

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  - $\longrightarrow$  most symmetries of H are not symmetries of  $\mathcal{T}$
- General proof (for any n): look at the commutative subgroups

• We already know a symmetry that twists the TPS: **the time evolution** (*Heisenberg picture*)

**Proposition.** The time evolution unitaries  $(e^{-it\hat{H}})_{t\in\mathbb{R}}$  are all symmetries of  $\mathcal{T}$  if and only if  $\hat{H}$  is 1-local with respect to  $\mathcal{T}$ .

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- Proof (sketch)
  - 1. think in Heisenberg picture

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  - 4. t  $\rightarrow$  S[  $\rho_i(t)$  ] takes an uncountable number of values, so  $(e^{-iHt})_t \cdot \mathcal{T}_0$  sweeps an continuous orbit of TPSs

# III. HOW STRUCTURES ACQUIRE THEIR IDENTITY

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C. Cao and S. M. Carroll, "Bulk entanglement gravity without a boundary: Towards finding Einstein's equation in Hilbert space", Physical Review D (2018)

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  - the relevant notion of uniqueness in physics is relational

NB: also in maths!

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**Proposition.** Let  $(\Pi_i^{\hat{H}})_{1 \leq i \leq n}$  denote the eigenprojectors of  $\hat{H}$  and  $\Lambda = (\lambda_i)_{1 \leq i \leq n}$  a family of non-negative real numbers such that  $\sum_i \lambda_i^2 = 1$ . The set:

$$\mathcal{K}_{HSF(\sigma,\Lambda)} = \{(\hat{H}, |\psi\rangle) \mid \operatorname{Spec}(\hat{H}) = \sigma \ and \ \forall i, \ \|\Pi_i^{\hat{H}} |\psi\rangle\|^2 = \lambda_i\}$$

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• In fact, the unitary group acts **transitively and freely** on this kind.

We can now formalize the intuition of III.1.

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**Definition** (Relational uniqueness). Let  $\mathcal{K}_0$  and  $\mathcal{K}_e$  be two determined kinds, and P a unitary-invariant property on the kind  $\mathcal{K}_0 \times \mathcal{K}_e$ . We say that P determines the product kind  $\mathcal{K}_0 \times \mathcal{K}_e$  if  $\{(\mathcal{S}_0, \mathcal{S}_e) \in \mathcal{K}_0 \times \mathcal{K}_e \mid P(\mathcal{S}_0, \mathcal{S}_e)\}$  is non-empty and if it is a determined kind, said differently:

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  - recall the proof of II.4. (uncountable infinity of TPSs)

#### Key result

**Proposition.** Let  $K_0$  and  $K_e$  be two determined kinds, and  $S_0 \in K_0$ . The following are equivalent:

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• The big question: is it possible to determine the kind  $\mathcal{K}_{HSF(\sigma, \Lambda)}$  x  $\mathcal{K}_{TPS(n; d1 ... dn)}$ ?

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P selects a unique triplet (H,  $|\psi\rangle$ , T) up to global unitary equivalence

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  - 6. Conclude by the following lemma:

**Lemma.** An operator  $A \in \mathcal{L}(\mathcal{H}_1 \otimes \cdots \otimes \mathcal{H}_n)$  is a product of single-site operators  $A_1 \otimes \cdots \otimes A_n$  if and only if A maps pure tensors to pure tensors.

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- Selecting an emergent structure *via* a physical principle has proved successful. Examples: Einstein's tensor [Lovelock 1971], quantum fields, Schrödinger equation, bosonic and fermonic statistics [Mekonnen *et al.*, 2025]...)

F. Klein, "Vergleichende Betrachtungen über neuere geometrische Forschungen", Mathematische Annalen (1893)

D. Lovelock, "The Einstein tensor and its generalizations", Journal of Mathematical Physics (1971)

M. Mekonnen, T. D. Galley, and M. P. Müller, "Invariance under quantum permutations rules out parastatistics", arXiv preprint (2025)

# Thank you for your attention!

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